



Heriot-Watt University  
Research Gateway

## Decomposition for coherence current of mutual coherence function for stochastic optical field

### Citation for published version:

Zhu, X, Beaumont, J, Chen, B, Takeda, M & Wang, W 2020, Decomposition for coherence current of mutual coherence function for stochastic optical field. in X Luo, Y Jiang, J Lu & D Liu (eds), *AOPC 2020: Optical Sensing and Imaging Technology.*, 1156705, Proceedings of SPIE, vol. 11567, SPIE, 2020 Applied Optics and Photonics China, Xiamen, China, 25/08/20. <https://doi.org/10.1117/12.2572765>

### Digital Object Identifier (DOI):

[10.1117/12.2572765](https://doi.org/10.1117/12.2572765)

### Link:

[Link to publication record in Heriot-Watt Research Portal](#)

### Document Version:

Publisher's PDF, also known as Version of record

### Published In:

AOPC 2020

### Publisher Rights Statement:

Copyright 2020 Society of PhotoOptical Instrumentation Engineers (SPIE). One print or electronic copy may be made for personal use only. Systematic reproduction and distribution, duplication of any material in this publication for a fee or for commercial purposes, and modification of the contents of the publication are prohibited.

Proceedings Volume 11567, AOPC 2020: Optical Sensing and Imaging Technology; 1156705 (2020)  
<https://doi.org/10.1117/12.2572765>

### General rights

Copyright for the publications made accessible via Heriot-Watt Research Portal is retained by the author(s) and / or other copyright owners and it is a condition of accessing these publications that users recognise and abide by the legal requirements associated with these rights.

### Take down policy

Heriot-Watt University has made every reasonable effort to ensure that the content in Heriot-Watt Research Portal complies with UK legislation. If you believe that the public display of this file breaches copyright please contact [open.access@hw.ac.uk](mailto:open.access@hw.ac.uk) providing details, and we will remove access to the work immediately and investigate your claim.

# PROCEEDINGS OF SPIE

[SPIDigitalLibrary.org/conference-proceedings-of-spie](https://SPIDigitalLibrary.org/conference-proceedings-of-spie)

## Decomposition for coherence current of mutual coherence function for stochastic optical field

Zhu, Xueliang, Beaumont, Jacob, Chen, Baixin, Takeda, Mitsuo, Wang, Wei

Xueliang Zhu, Jacob Beaumont, Baixin Chen, Mitsuo Takeda, Wei Wang, "Decomposition for coherence current of mutual coherence function for stochastic optical field," Proc. SPIE 11567, AOPC 2020: Optical Sensing and Imaging Technology, 1156705 (5 November 2020); doi: 10.1117/12.2572765

**SPIE.**

Event: Applied Optics and Photonics China (AOPC 2020), 2020, Beijing, China

# Decomposition for Coherence Current of Mutual Coherence Function for Stochastic Optical Field

Xueliang Zhu<sup>a,b</sup>, Jacob Beaumont<sup>a</sup>, Baixin Chen<sup>a</sup>, Mitsuo Takeda<sup>b,c</sup> and Wei Wang<sup>a,b\*</sup>

<sup>a</sup> Institute of Photonics and Quantum Sciences, School of Engineering and Physical Science, Heriot-Watt University, Edinburgh, EH14 4AS, United Kingdom;

<sup>b</sup> international Center for Optical Research and Education (iCORE), Xi'an Technological University, Xi'an, 710032, P. R. China;

<sup>c</sup> Center for Optical Research and Education (CORE), Utsunomiya University, 7-1-2 Yoto, Utsunomiya, Tochigi 321-8585, Japan

## ABSTRACT

In analogy to the velocity decomposition in continuum mechanics, we introduce two new tensors, referred to as the deformation-rate tensor and rotation-rate tensor, to the optical coherence theory for stochastic optical field, and decompose a coherence current vector into its translation, rotation and deformation components to study the optical coherence dynamics. To investigate the optical coherence propagation and evolution, we have conducted an experiment with results given to demonstrate the two newly introduced tensors.

**Keywords:** Coherence Current, Optical Coherence Function, Stochastic Optical Field

## 1. INTRODUCTION

Optical fields are inherently of a statistical nature and the cross-correlation between the fluctuating fields known as the coherence function is a quantity of great importance in understanding the statistics of stochastic optical field. [1-2] During the recent years, a new concept of coherence current has been introduced and observed when a generic coherence vortex was experimentally observed. [3-4] Just as the unique role of Poynting vector in Electrodynamics, the coherence current plays a critical role to describe the propagation of optical coherence, where its magnitude is proportional to the fringe contrast and its vectorial direction points along the propagation of optical coherence.

In continuum mechanics, the strain-rate tensor or rate-of-strain tensor is a physical quantity that describes the rate of change of the deformation of a material. It can be defined as the derivative of the strain tensor with respect to time, or as the symmetric component of the gradient (derivative with respect to position) of the flow velocity. In fluid mechanics it also can be described as the velocity gradient, a measure of how the velocity of a fluid changes between different points within the fluid.[5] The concept has implications in a variety of areas of physics and engineering, including magnetohydrodynamics, mining and water treatment. Since the strain rate tensor is a purely kinematic concept that describes the macroscopic motion of the material, therefore, it does not depend on the nature of the material, or on the forces and stresses that may be acting on it; and it applies to any continuous medium, whether solid, liquid or gas.

In this paper, we will follow the approach of velocity decomposition developed for fluid dynamics, and decompose the coherence current vector for optical coherence function of stochastic optical field. After introduction of two new tensors, referred to as the deformation-rate tensor and rotation-rate tensor, we are able to decompose a coherence current vector into its translation, rotation and deformation components to study the optical coherence dynamics. Experiments have been conducted to visualize these two newly introduced tensors to investigate the optical coherence propagation and evolution.

---

\*[w.wang@hw.ac.uk](mailto:w.wang@hw.ac.uk) Tel: +44 (0) 131 451 3141; Fax: +44 (0) 131 451 3129

## 2. DECOMPOSITION OF COHERENCE CURRENT

To understand the kinematics of optical coherence, we can start our analysis from the coherence current given by  $\mathbf{T}(\mathbf{r}_1, \mathbf{r}_2) = -j\alpha\bar{k}[J^*\nabla_1 J - J\nabla_1 J^*]$ , where  $\alpha$  is a positive constant,  $\bar{k} = 2\pi\bar{\nu}/c$  is the mean value of wave number,  $c$  is the constant light speed,  $J(\mathbf{r}_1, \mathbf{r}_2) = \langle u^*(\mathbf{r}_1)u(\mathbf{r}_2) \rangle$  is the mutual intensity of optical field with  $\langle \cdot \rangle$  indicating an ensemble average [4-5]. Without loss of generality, we will restrict our discussions to the mutual coherence function for variation of location for the  $\mathbf{r}_1$  by keeping  $\mathbf{r}_2$  fixed. Similar to the well-known Poynting vector in an electromagnetic field, the coherence current vector  $\mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)$ , proportional to the linear coherence momentum, represents the directional optical coherence flux with its magnitude indicating the fringe contrast and its vectorial direction pointing in the direction of propagation. The coherence current displaced from  $\mathbf{r}_1$  by a small vector  $d\mathbf{r}_1$ , can be approximated after a Taylor series expansion:

$$\mathbf{T}(\mathbf{r}_1 + d\mathbf{r}_1, \mathbf{r}_2) \approx \mathbf{T}(\mathbf{r}_1, \mathbf{r}_2) + \nabla_1 \mathbf{T}(\mathbf{r}_1, \mathbf{r}_2) \cdot d\mathbf{r}_1, \quad (1)$$

where  $\nabla_1 \mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)$ , the gradient of the coherence current is a  $3 \times 3$  tensor which could be understood as a linear map that takes a displacement vector  $d\mathbf{r}_1$  to the corresponding change in the coherence current. Note the fact that any matrix can be decomposed into the sum of a symmetric matrix and an anti-symmetric matrix. Applying this to  $\nabla_1 \mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)$  and following the same operation used in continuum mechanics, we arrive at a sum of symmetric and antisymmetric components. Thus, the coherence current may be rewritten as:

$$\mathbf{T}(\mathbf{r}_1 + d\mathbf{r}_1, \mathbf{r}_2) = \mathbf{T}(\mathbf{r}_1, \mathbf{r}_2) + \boldsymbol{\Omega}(\mathbf{r}_1, \mathbf{r}_2) \times d\mathbf{r}_1 + \boldsymbol{\varepsilon}(\mathbf{r}_1, \mathbf{r}_2) \cdot d\mathbf{r}_1, \quad (2)$$

where ' $\times$ ' and ' $\cdot$ ' are tensor's cross product and dot product, respectively,  $\boldsymbol{\Omega}$  and  $\boldsymbol{\varepsilon}$  are the *rate of coherence rotation* tensor and the *rate of coherence deformation* tensor given by

$$\boldsymbol{\Omega}(\mathbf{r}_1, \mathbf{r}_2) = \begin{bmatrix} 0 & -\Omega_z(\mathbf{r}_1, \mathbf{r}_2) & \Omega_y(\mathbf{r}_1, \mathbf{r}_2) \\ \Omega_z(\mathbf{r}_1, \mathbf{r}_2) & 0 & \Omega_x(\mathbf{r}_1, \mathbf{r}_2) \\ -\Omega_y(\mathbf{r}_1, \mathbf{r}_2) & -\Omega_x(\mathbf{r}_1, \mathbf{r}_2) & 0 \end{bmatrix} = \begin{bmatrix} 0 & -0.5\left(\frac{\partial T_y}{\partial x_1} - \frac{\partial T_x}{\partial y_1}\right) & 0.5\left(\frac{\partial T_x}{\partial z_1} - \frac{\partial T_z}{\partial x_1}\right) \\ 0.5\left(\frac{\partial T_y}{\partial x_1} - \frac{\partial T_x}{\partial y_1}\right) & 0 & 0.5\left(\frac{\partial T_z}{\partial y_1} - \frac{\partial T_y}{\partial z_1}\right) \\ -0.5\left(\frac{\partial T_x}{\partial z_1} - \frac{\partial T_z}{\partial x_1}\right) & -0.5\left(\frac{\partial T_z}{\partial y_1} - \frac{\partial T_y}{\partial z_1}\right) & 0 \end{bmatrix}, \quad (3)$$

and

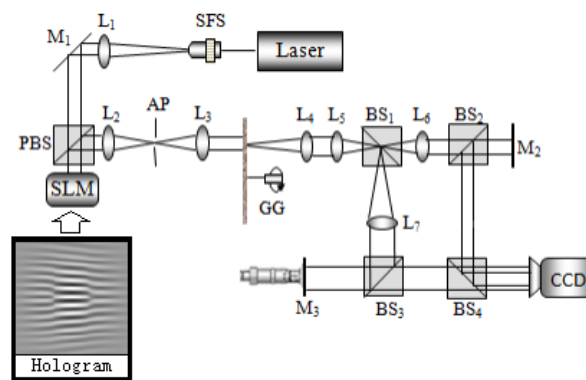
$$\boldsymbol{\varepsilon}(\mathbf{r}_1, \mathbf{r}_2) = \begin{bmatrix} \varepsilon_{xx}(\mathbf{r}_1, \mathbf{r}_2) & \varepsilon_{xy}(\mathbf{r}_1, \mathbf{r}_2) & \varepsilon_{xz}(\mathbf{r}_1, \mathbf{r}_2) \\ \varepsilon_{yx}(\mathbf{r}_1, \mathbf{r}_2) & \varepsilon_{yy}(\mathbf{r}_1, \mathbf{r}_2) & \varepsilon_{yz}(\mathbf{r}_1, \mathbf{r}_2) \\ \varepsilon_{zx}(\mathbf{r}_1, \mathbf{r}_2) & \varepsilon_{zy}(\mathbf{r}_1, \mathbf{r}_2) & \varepsilon_{zz}(\mathbf{r}_1, \mathbf{r}_2) \end{bmatrix} = \begin{bmatrix} \frac{\partial T_x}{\partial x_1} & 0.5\left(\frac{\partial T_y}{\partial x_1} + \frac{\partial T_x}{\partial y_1}\right) & 0.5\left(\frac{\partial T_x}{\partial z_1} + \frac{\partial T_z}{\partial x_1}\right) \\ 0.5\left(\frac{\partial T_y}{\partial x_1} + \frac{\partial T_x}{\partial y_1}\right) & \frac{\partial T_y}{\partial y_1} & 0.5\left(\frac{\partial T_z}{\partial y_1} + \frac{\partial T_y}{\partial z_1}\right) \\ 0.5\left(\frac{\partial T_x}{\partial z_1} + \frac{\partial T_z}{\partial x_1}\right) & 0.5\left(\frac{\partial T_z}{\partial y_1} + \frac{\partial T_y}{\partial z_1}\right) & \frac{\partial T_z}{\partial z_1} \end{bmatrix}. \quad (4)$$

The antisymmetric tensor  $\boldsymbol{\Omega}(\mathbf{r}_1, \mathbf{r}_2)$  represents a rigid body rotation of the coherence current along the running point  $\mathbf{r}_1$  where the corresponding coherence vorticity vector is given by  $(\Omega_x, \Omega_y, \Omega_z)^t = \nabla_1 \times \mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)/2$  with a superscript ' $t$ ' indicating a matrix transpose. The coherence vorticity is a key characteristic of coherence current and have been used to explain the coherence vortex [4-5]. The symmetric coherence deformation rate tensor  $\boldsymbol{\varepsilon}(\mathbf{r}_1, \mathbf{r}_2)$  describes optical coherence expansion or contraction  $\varepsilon_{mm}, (m=n)$  and optical coherence shear  $\varepsilon_{mn}, (m \neq n)$  resulting from optical coherence gradients in the inner of the volume. Therefore, we can refer to the diagonal components in this tensor as the normal strain rates of optical coherence and the off-diagonal components as shear strain rates of optical coherence. From

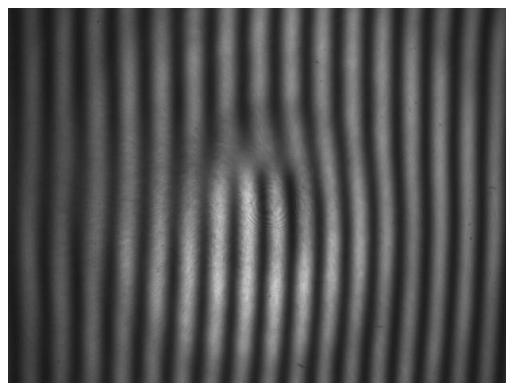
Eq. (2)-(4), we can see that, to the first order in the linear dimensions of a small region surrounding the position  $\mathbf{r}_1$ , the coherence current vector  $\mathbf{T}$  stemming from the correlation of optical fields at the point  $\mathbf{r}_1 + d\mathbf{r}_1$  and  $\mathbf{r}_2$  consists, in effect, of the superposition of a uniform translation of coherence current  $\mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)$ , a pure deforming motion characterized by the rate of optical coherence deformation tensor  $\mathbf{\epsilon}(\mathbf{r}_1, \mathbf{r}_2)$  which itself may be decomposed into an isotropic expansion or contraction and a shearing motion without change of its volume, and a rigid-body rotation with angular velocity  $(\Omega_x, \Omega_y, \Omega_z)^t$ .

### 3. EXPERIMENTAL DEMONSTRATION

To study the evolution of the coherence currents associated with a coherence vortex, an experiment has been conducted to synthesize a phase singularity in the optical coherence function. Fig.2(a) shows the experimental geometry for coherence synthesis and visualization based on the principle of coherence holography [7]. A computer-generated hologram encoded with vortex information is projected onto a rotating ground glass to serve as a spatially incoherent light source. For coherence visualization, two telescope systems with different magnification have been used for two arms of the Michelson interferometer to introduce the radial shearing. From the recorded interferogram as shown in Fig.1(b), we have reconstructed the complex-valued mutual intensity and measured the visibility of the fringes by using the Fourier transform method [8]. By calculation of coherence current from its definition, we obtained the distribution of the coherence current vector and conducted its decomposition from the two newly introduced tensors based on their definition to study the evolution of the optical coherence for stochastic optical field.



(a)



(b)

Fig.1. Experimental setup of coherence holography (a); and the recorded interferogram (b).

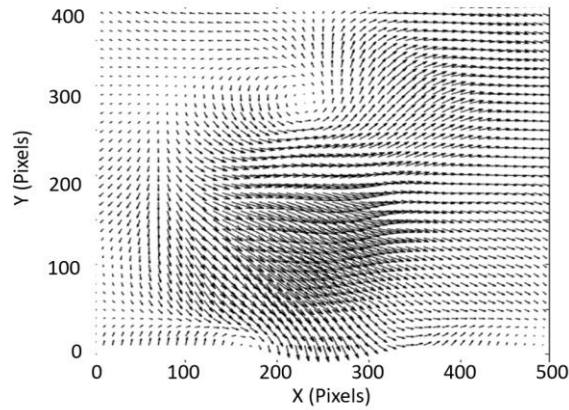


Fig.2 Quiver plot of the coherence currents reconstructed from the recorded interferogram

Figure 2 shows the quiver plot of the coherence current reconstructed from the recorded interferogram as shown in Fig. 1(b), where the length of each arrow indicates the magnitude of local coherence current with its arrow direction pointing out the corresponding orientation of coherence current vector. As expected, the coherence current circulates around the coherence vortex located at the centers of fork-like fringe pattern.

Since we are only carrying out a 2-dimensional analysis of the coherence current from the recorded interferogram, the rate of rotation tensor  $\mathbf{\Omega}$  and rate of deformation tensor  $\mathbf{\epsilon}$  has been simplified by eliminating some tensor components associated with the  $\hat{z}$  coordinate. Following such a simplification, the tensor components  $\Omega_z, \epsilon_{xx}, \epsilon_{yy}$  and  $\epsilon_{xy}(=\epsilon_{yx})$  remain and can be reconstructed from the numerical calculations based on their definitions in Eqs. (3) and (4).

Figure 3 shows the distribution of  $\Omega_z$  indicating the distribution of the vorticity as the rate of rotation around the  $\hat{z}$  axis. As expected, a large value of  $\Omega_z$  can be observed at the location where a coherence vortex exists. When the value of  $\Omega_z$  becomes negative, the local direction for the rate of rotation changes its direction.

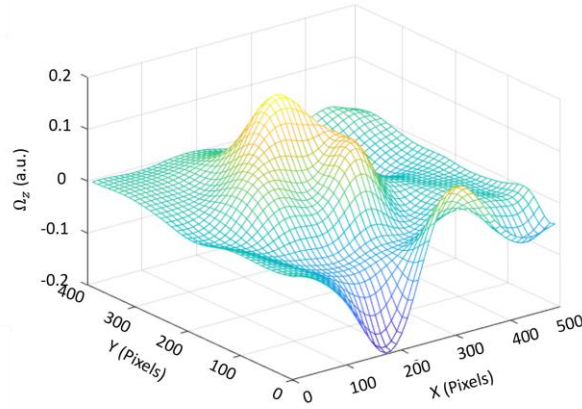
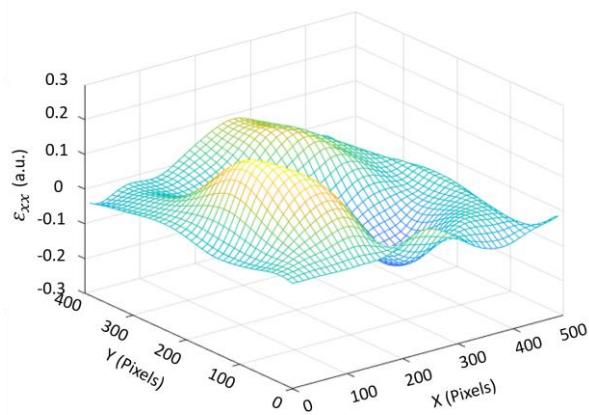
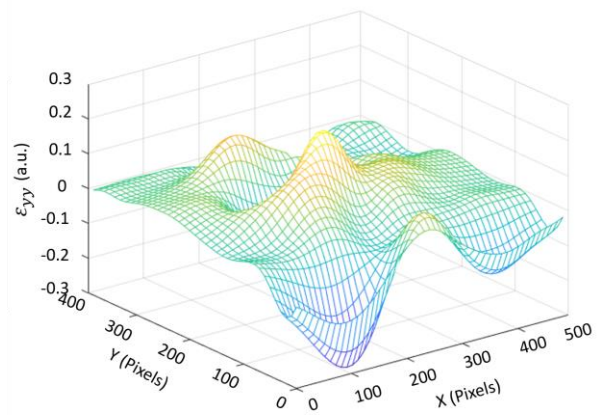


Fig. 3 Distribution of the tensor component  $\Omega_z$  for the rate of coherence rotation tensor



(a)



(b)

Fig. 4 Distribution of the tensor components  $\varepsilon_{xx}$  and  $\varepsilon_{yy}$  for the rate of coherence deformation tensor

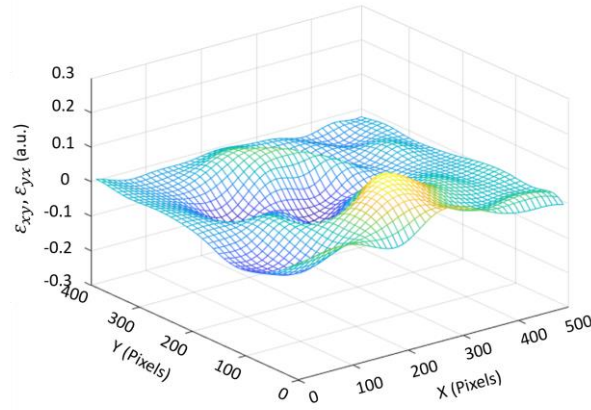


Fig. 5 Distribution of the tensor components  $\varepsilon_{xy}$  and  $\varepsilon_{yx}$  for the rate of coherence deformation tensor

Figure 4(a) and (b) give the distributions of the normal rates of deformation for the coherence current  $\mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)$  along the  $\hat{x}$  and  $\hat{y}$  directions, respectively, where the normal deformation rates for this example can be observed and analyzed from these diagonal tensor components  $\varepsilon_{xx}$  and  $\varepsilon_{yy}$ . The corresponding shearing deformation rates of  $\mathbf{T}(\mathbf{r}_1, \mathbf{r}_2)$ , i.e.  $\varepsilon_{xy} (= \varepsilon_{yx})$  have been given in Fig. 5 indicating a distribution of change rate in shape for the coherence current.

#### 4. CONCLUSIONS

In summary, we have adopted the theorem developed for fluid dynamics to study the dynamics of the optical coherence for stochastic optical field. By introducing the deformation-rate tensor and rotation-rate tensor to statistical optics, we have investigated the evolution of the coherence current by its decomposition into a translation, rotation and deformation contributions. Experiments have been conducted to demonstrate these two newly introduced tensors.

## REFERENCES

- [1] M. Born, E. Wolf, [Principles of Optics], Cambridge University Press, (1980).
- [2] J. W. Goodman, [Statistical Optics], 2<sup>nd</sup> edition, Wiley, Colorado (2015).
- [3] H. F. Schouten, G. V. Bogatyryova, C. V. Fel'de, P. V. Polyanskii, S. A. Ponomarenko, M. S. Soskin, and E. Wolf, "Partially coherent vortex beams with a separable phase," Opt. Lett. 28, 878 (2003).
- [4] W. Wang, Z. Duan, S. G. Hanson, Y. Miyamoto, and M. Takeda, "Experimental Study of Coherence Vortices: Local Properties of Phase Singularities in a Spatial Coherence Function," Phys. Rev. Lett. 96, 073902 ,(2006).
- [5] W. Wang and M. Takeda, "Coherence Current, Coherence Vortex, and the Conservation Law of Coherence," Phys Rev. Lett, 96(22): 223904(2006).
- [6] G. K. Batchelor, An Introduction to Fluid Dynamics, Chapter 2 (Cambridge University Press, Cambridge, 1967).
- [7] M. Takeda, W. Wang, Z. Duan, Y. Miyamoto, "Coherence Holography," Opt. Express 23, 9629-9635 (2005).
- [8] M. Takeda, H. Ina and S. Kobayashi, "Fourier-transform method of fringe-pattern analysis for computer-based topography and interferometry," JOSA 72, 156-160 (1982).